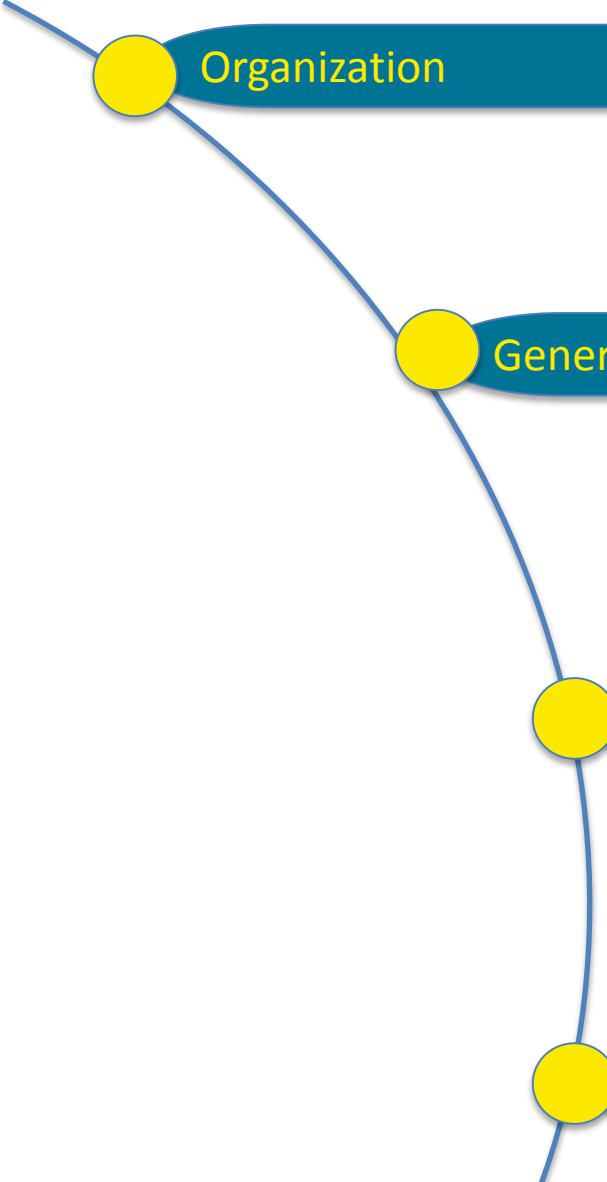




Organization

General objectives

A variety of platforms

Highlights

| Quantum sensors and Metrology    |              |
|----------------------------------|--------------|
| <b>Franck Pereira dos Santos</b> | <b>SYRTE</b> |
| <b>Ivan Favero</b>               | <b>MPQ</b>   |
| Sophie Djordjevic                | LNE          |
| Fabien Bretenaker                | LAC          |
| Benoit Darquié                   | LPL          |

**Number of groups involved : 32**

Do not hesitate to contact us  
if you have any question related to the calls for projects, the evaluation ...

## Scope of the area

- Study quantum properties of matter and light to realize sensors
- Measure physical quantities to probe the environment
- Realize standards for these measurements

## How to ?

By exploiting the principles of quantum mechanics

Central to our field is the quantum superposition, the question of the coherence and the use of non classical states and entanglement as a resource

### Goals :

Beat the classical limits, reach the quantum limit, go beyond ...

Quantum limit ? Could be defined as « **a limit for noise levels set by quantum mechanics** »

### Technological aspects:

How to use these sensors? How to deploy them ? What level of integration can be reached ?  
For which applications?

## Different platforms

**Atoms/molecules/artificial atoms/photons/optomechanical devices/mesoscopic systems**  
**Different levels of maturity, different enabling technologies, different measurement principles**

## A few examples

### Atomic sensors:

MW clocks and inertial sensors in differential mode at the QPN

Optical clocks/inertial sensors still limited by « classical » noise (laser / vibration noise)

Demonstration of optical magnetometers below the SQL

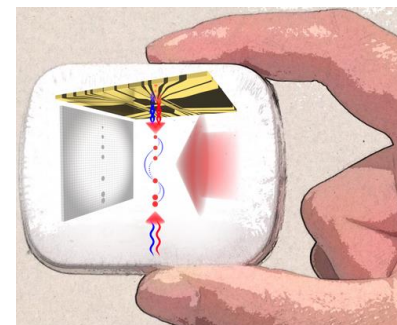
**Maturity:** technology transfer



**Deployment:** underway



**Integration:**  
still requires efforts

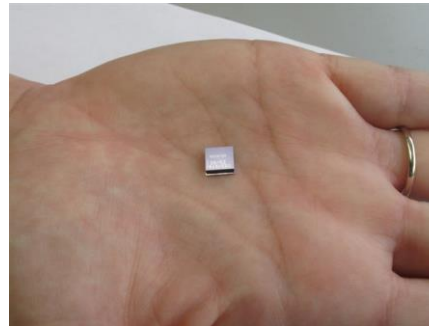
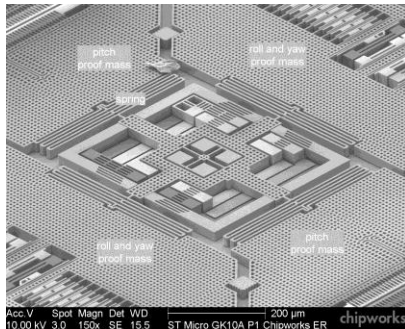


**A few examples**

**Optomechanical sensors:** SQL reached recently (Peterson et al., PRL 116, 063601 (2016))  
With a very strong optomechanical coupling, fluctuations are imposed by the quantum fluctuations of light (rather than by thermal noise)  
Key: very low T ... but should be possible at high temperature if increasing the coupling.

**Integration:** OK (nano-objets).

**Deployment:** targeted / projects ongoing



**Artificial atoms:** N/V centers in diamond, vacancies in SiC, dopants in crystals ...

**Temperature sensors, magnetometers, electrometers, strain sensors at the nanometer scale**

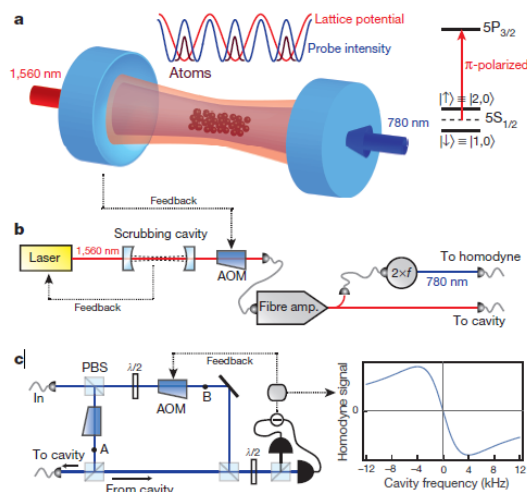
Detection of spins by ESR with sensitivity limited by quantum fluctuations of the electromagnetic field instead of thermal or technical noise. Requires amplifiers at very low T.

Challenge: « high temperature » materials

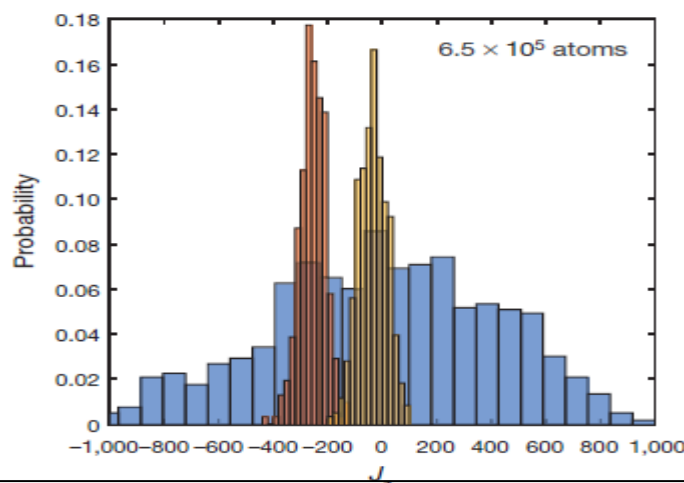
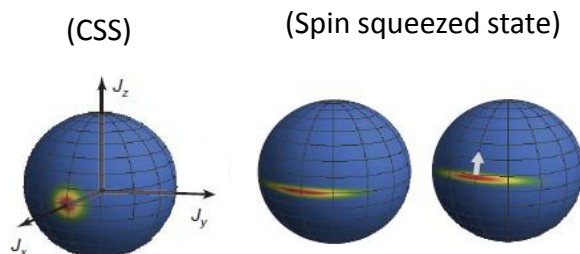
# Squeezing the collective spin of an ensemble of $5 \cdot 10^5$ atoms via an optical-cavity-based measurement

(Hosten et al., Nature (2016) « Measurement noise 100 times lower than the quantum-projection limit using entangled atoms »)

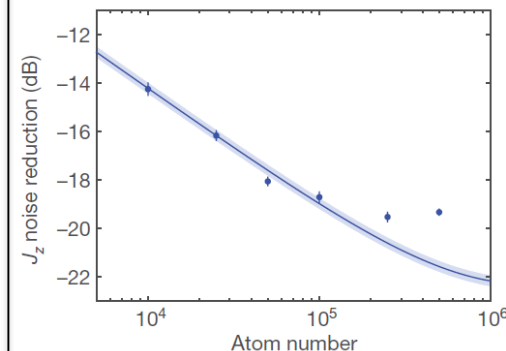
**Preparation of a spin squeezed state through a collective population difference measurement**



**A  $J_z$  measurement projects the quantum state into one with a narrower distribution of  $J_z$  than a coherent spin state (CSS)**



**Demonstrated: 20 dB squeezing**



**Proof of principle experiment in clock configuration : 10.5 dB gain**

**Remains to be implemented in a state of the art interferometer (clock, inertial sensor)**

**Similar setups currently developed/under development on clocks at SYRTE**

## A sensitive electrometer based on a Ry atom in a Schrodinger-cat state

A. Facon, E.-K. Dietsche, D. Grosso, S. Haroche, J.-M. Raimond, M. Brune & S. Gleyzes, *Nature* 18335 (2016)

### Measurement of the electric field using a large angular momentum (quantum number $J \approx 25$ ) with a single atom in a high energy Ry state

Atoms prepared in  $|J, J\rangle$

Standard Ramsey interferometer

$$\phi = (\omega(F) - \omega_{\text{rf}})\tau$$

Probability to find the atoms in  $|J, J\rangle$

$$P(\varphi_{\text{rf}}) = |\langle \theta, \varphi_{\text{rf}} | \theta, \phi \rangle|^2 = \exp(-J \sin^2 \theta (\phi - \varphi_{\text{rf}})^2 / 2)$$



$$\langle J, J | \psi_f \rangle = |\langle J, J | \psi_f \rangle| e^{-i\Phi} \quad \Phi = J \left[ \phi - \varphi_{\text{rf}} - 2 \arctan \left[ \cos \theta \tan \left( \frac{\phi - \varphi_{\text{rf}}}{2} \right) \right] \right]$$

Add MW pulses, coupling to reference state

$$(|J, J\rangle + |R\rangle) / \sqrt{2}$$

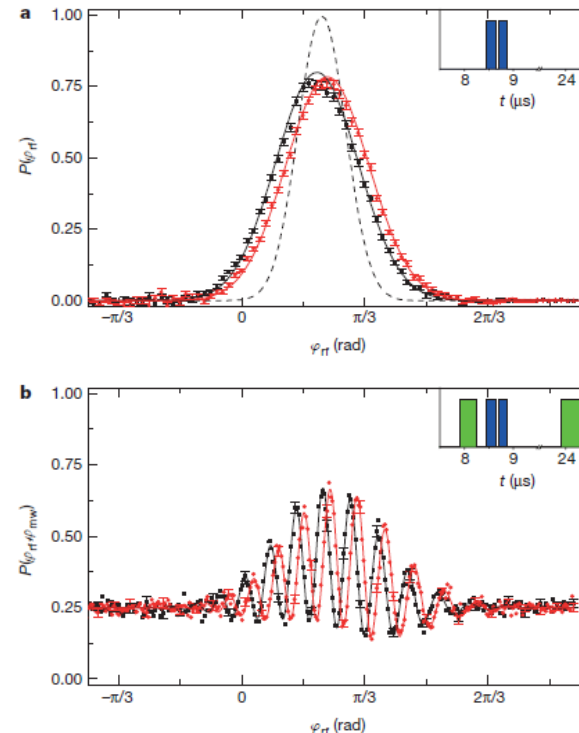
Make them interfere

Probability to find the atoms in  $|J, J\rangle$

$$P(\varphi_{\text{rf}}, \varphi_{\text{mw}}) = \frac{1}{4} + \frac{1}{4}P(\varphi_{\text{rf}}) + \frac{1}{2}\sqrt{P(\varphi_{\text{rf}})} \cos(\Phi - \varphi_{\text{mw}})$$

Finally, sensitivity outperforms the SQL by a factor of up to 2

RF pulses at  $\omega_{\text{rf}}$



$$\sigma_{F, \text{SQL}}^1 = \frac{1}{\tau \sqrt{2J}} \left| \frac{\partial \omega}{\partial F} \right|^{-1}$$

**$1/J^{1/2}$  scaling  
(CSS)**

$$\sigma_F^{1, \text{th}} = \frac{1}{\tau J (1 - \cos \theta)} \left| \frac{\partial \omega}{\partial F} \right|^{-1}$$

**$1/J$  scaling**